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A little tedious algebra can show that the transfer function above has two poles and a zero

$$H(0) \approx -g_m(R \parallel r_o) \qquad \frac{1}{\tau_3} = \frac{g_m}{C_{gd}}$$

1) Assuming $g_m R_s >> 1$, these poles are:

$$\frac{1}{\tau_1} \approx \frac{g_m}{C_{gs}} + \frac{1}{C_{gd}(R \parallel r_o)} \qquad \qquad \frac{1}{\tau_2} \approx \frac{1}{\left[C_{gs} + C_{gd}(g_m(R \parallel r_o))\right] R_s}$$

This pole will likely determine the smallest frequency at which the total gain rolls over

2) Assuming $g_m R_s \ll 1$, these poles are:

 $\frac{1}{\tau_1} \approx \frac{1}{R_s C_{gs}}$

$$\frac{1}{\tau_2} \approx \frac{1}{C_{gd}(R \parallel r_o)}$$

This pole will likely determine the smallest frequency at which the total gain rolls over

The Common Source Amplifier: The Miller Approximation

If the poles and zeros in $A_v(\omega)$ are at a higher frequency than the poles and zeros associated with the transfer function $v_{in}(\omega)/v_s(\omega)$ (which would be the case if $g_m R_s >> 1$) then one may approximate the open circuit gain $A_v(\omega)$ by its low frequency value:

$$A_{v}(\omega) = \frac{v_{out}(\omega)}{v_{in}(\omega)} = -\frac{g_{m}R - j\omega RC_{gd}}{1 + g_{o}R + j\omega RC_{gd}} \approx -g_{m}(R || r_{o})$$

Miller approximation

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And then:

$$H(\omega) = \frac{v_{out}(\omega)}{v_s(\omega)} = \frac{v_{out}(\omega)}{v_{in}(\omega)} \frac{v_{in}(\omega)}{v_s(\omega)} = A_v(\omega) \frac{Z_{in}(\omega)}{R_s + Z_{in}(\omega)}$$
$$\approx \frac{-g_m(R || r_o)}{1 + j\omega [C_{gs} + C_{gd}(1 + g_m(R || r_o))] R_s}$$

Looking in from the input terminal the capacitance C_{gd} seems to be magnified by a factor proportional to the low frequency open circuit voltage gain of the amplifier!!

And now the single pole is at:

$$\frac{1}{\tau} \approx \frac{1}{\left[C_{gs} + C_{gd}(1 + g_m(R || r_o))\right] R_s}$$



















